## How Einstein derived Planck's law

Consider a gas at temperature $T$. Let $N$ be the number of gas molecules and let $N_{n}$ be the number of molecules with energy $E_{n}$. By the Maxwell-Boltzmann distribution we have

$$
\begin{equation*}
\frac{N_{n}}{N}=p_{n} \exp \left(-\frac{E_{n}}{k T}\right) \tag{1}
\end{equation*}
$$

Coefficient $p_{n}$ is a statistical weighting factor that does not depend on $T$.
Let us now consider the processes by which an atom or molecule transitions between energy levels. The processes are absorption, induced emission, and spontaneous emission. Let $E_{m}$ and $E_{n}$ be energy levels such that $E_{m}>E_{n}$. Let $N_{m \rightarrow n} / \Delta t$ be the number of atoms or molecules that transition from energy level $E_{m}$ to $E_{n}$ in time $\Delta t$. Finally, let $B_{n m}, B_{m n}$, and $A_{m n}$ be coefficients such that

Absorption and induced emission are proportional to radiant energy density $\rho$. The $A$ and $B$ coefficients do not depend on $T$.

At equilibrium, the transition rates are equal.

$$
\frac{N_{n \rightarrow m}}{\Delta t}=\frac{N_{m \rightarrow n}}{\Delta t}
$$

Hence

$$
\begin{aligned}
& \rho B_{n m} N_{n}=\rho B_{m n} N_{m}+A_{m n} N_{m} \\
& \text { absorption } \begin{array}{c}
\text { induced } \\
\text { emission }
\end{array} \begin{array}{c}
\text { spontaneous } \\
\text { emission }
\end{array}
\end{aligned}
$$

Divide through by $N$.

$$
\rho B_{n m} \frac{N_{n}}{N}=\underset{\text { absorption }}{\rho B_{m n} \frac{N_{m}}{N}+A_{m n} \frac{N_{m}}{N}} \underset{\substack{\text { induced } \\ \text { emission }}}{\substack{\text { spontaneous } \\ \text { emission }}}
$$

Then by the Maxwell-Boltzmann distribution (1) we have

$$
\begin{equation*}
\rho B_{n m} p_{n} \exp \left(-\frac{E_{n}}{k T}\right)=\rho B_{m n} p_{m} \exp \left(-\frac{E_{m}}{k T}\right)+A_{m n} p_{m} \exp \left(-\frac{E_{m}}{k T}\right) \tag{2}
\end{equation*}
$$

Multiply both sides by $\exp \left(E_{m} / k T\right)$.

$$
\rho B_{n m} p_{n} \exp \left(\frac{E_{m}-E_{n}}{k T}\right)=\underset{\substack{\text { induced } \\ \text { emission }}}{\rho B_{m n} p_{m}}+\underset{\substack{\text { spontaneous } \\ \text { emission }}}{A_{m n} p_{m}}
$$

Note that for increasing $T$ we have

$$
\lim _{T \rightarrow \infty} \exp \left(\frac{E_{m}-E_{n}}{k T}\right)=1
$$

It follows that for $T \rightarrow \infty$ the equilibrium formula is

$$
\rho B_{n m} p_{n}=\rho B_{m n} p_{m}+A_{m n} p_{m}
$$

Divide through by $\rho$.

$$
B_{n m} p_{n}=B_{m n} p_{m}+\frac{A_{m n} p_{m}}{\rho}
$$

Energy density $\rho$ increases with temperature $T$ hence $A_{m n} p_{m} / \rho$ vanishes for $T \rightarrow \infty$ leaving

$$
\begin{equation*}
B_{n m} p_{n}=B_{m n} p_{m} \tag{3}
\end{equation*}
$$

Einstein reasoned that equation (3) is true in general based on the assumption that the factors involved do not depend on $T$. Hence we can substitute $B_{m n} p_{m}$ into the absorption term and write

$$
\rho B_{m n} p_{m} \exp \left(\frac{E_{m}-E_{n}}{k T}\right)=\underset{\substack{\text { induced } \\ \text { emission }}}{\rho B_{m n} p_{m}}+\underset{\substack{\text { spontaneous } \\ \text { emission }}}{A_{m n} p_{m}}
$$

Divide both sides by $B_{m n} p_{m}$.

$$
\rho \exp \left(\frac{E_{m}-E_{n}}{k T}\right)=\rho+\frac{A_{m n}}{B_{m n}}
$$

Rearrange terms.

$$
\rho \exp \left(\frac{E_{m}-E_{n}}{k T}\right)-\rho=\frac{A_{m n}}{B_{m n}}
$$

Factor out $\rho$.

$$
\rho\left(\exp \left(\frac{E_{m}-E_{n}}{k T}\right)-1\right)=\frac{A_{m n}}{B_{m n}}
$$

Solve for $\rho$.

$$
\rho=\frac{A_{m n}}{B_{m n}} \frac{1}{\exp \left(\frac{E_{m}-E_{n}}{k T}\right)-1}
$$

We now consider the case of large exponentials such that

$$
\exp \left(\frac{E_{m}-E_{n}}{k T}\right) \approx \exp \left(\frac{E_{m}-E_{n}}{k T}\right)-1
$$

Hence for large exponentials

$$
\rho \approx \frac{A_{m n}}{B_{m n}} \exp \left(-\frac{E_{m}-E_{n}}{k T}\right)
$$

By equivalence with Wien's law (which is accurate for large $\nu$ ) we have

$$
\rho=\frac{2 h \nu^{3}}{c^{2}} \exp \left(-\frac{h \nu}{k T}\right)
$$

Hence

$$
\begin{equation*}
\frac{A_{m n}}{B_{m n}}=\frac{2 h \nu^{3}}{c^{2}} \tag{4}
\end{equation*}
$$

and

$$
E_{m}-E_{n}=h \nu
$$

Then by substitution we obtain Planck's law.

$$
\begin{aligned}
\rho & =\frac{A_{m n}}{B_{m n}} \frac{1}{\exp \left(\frac{E_{m}-E_{n}}{k T}\right)-1} \\
& =\frac{2 h \nu^{3}}{c^{2}} \frac{1}{\exp \left(\frac{h \nu}{k T}\right)-1}
\end{aligned}
$$

Let us now consider the values of the $A$ and $B$ coefficients. The coefficient for spontaneous emission can be computed from quantum mechanics. For example, for hydrogen transition $2 p \rightarrow 1 s$ we have

$$
A_{21}=\frac{16 e^{8}}{6561 \varepsilon_{0}^{4} h^{4} c^{3} a_{0}}=6.26 \times 10^{8} \text { second }^{-1}
$$

The coefficient for induced emission can be obtained from equation (4).

$$
B_{m n}=\frac{c^{2}}{2 h \nu^{3}} A_{m n}
$$

The coefficient for absorption can be computed from equation (3).

$$
B_{n m}=\frac{p_{m}}{p_{n}} B_{m n}
$$

The ratio $p_{m} / p_{n}$ is equal to $g_{m} / g_{n}$ where $g$ is the multiplicity for quantum numbers $\ell$ and $m_{s}$.

$$
g=(2 \ell+1)\left(2 m_{s}+1\right)
$$

Hence for hydrogen $2 p \rightarrow 1 s$ we have

$$
\begin{array}{ll}
g_{1}=2 & \left(\ell=0, m_{s}=1 / 2\right) \\
g_{2}=6 & \left(\ell=1, m_{s}=1 / 2\right)
\end{array}
$$

(Recall that $\ell=0$ for orbital $s$ and $\ell=1$ for orbital $p$.)

