Atomic transitions 1

Let $\Psi(\mathbf{r},t)$ be the following linear combination of two wave functions where $c_a(t)$ and $c_b(t)$ are dimensionless time-dependent coefficients such that $|c_a(t)|^2+|c_b(t)|^2=1$.

$$\Psi(\mathbf{r},t)=c_a(t)\psi_a(\mathbf{r})\exp\left(-\tfrac{i}{\hbar}E_a t\right)+c_b(t)\psi_b(\mathbf{r})\exp\left(-\tfrac{i}{\hbar}E_b t\right)$$

Let the Hamiltonian be

$$H(\mathbf{r},t)=H_0(\mathbf{r})+H_1(\mathbf{r},t)$$

where

$$H_0\psi_a=E_a\psi_a,\quad H_0\psi_b=E_b\psi_b$$

We want to find solutions for $c_a(t)$ and $c_b(t)$. Start with the Schrödinger equation.

$$i\hbar\frac{\partial}{\partial t}\Psi=H\Psi$$

Evaluate the left side of the Schrödinger equation.

$$\begin{aligned}i\hbar\frac{\partial}{\partial t}\Psi=&\overbrace{E_a c_a(t)\psi_a(\mathbf{r})\exp\left(-\tfrac{i}{\hbar}E_a t\right)+E_b c_b(t)\psi_b(\mathbf{r})\exp\left(-\tfrac{i}{\hbar}E_b t\right)}^{\text{cancels with other side of Schrödinger equation}}\\&+i\hbar\dot{c}_a(t)\psi_a(\mathbf{r})\exp\left(-\tfrac{i}{\hbar}E_a t\right)+i\hbar\dot{c}_b(t)\psi_b(\mathbf{r})\exp\left(-\tfrac{i}{\hbar}E_b t\right)\end{aligned}$$

Evaluate the right side of the Schrödinger equation.

$$H\Psi=\overbrace{E_a c_a(t)\psi_a(\mathbf{r})\exp\left(-\tfrac{i}{\hbar}E_a t\right)+E_b c_b(t)\psi_b(\mathbf{r})\exp\left(-\tfrac{i}{\hbar}E_b t\right)}^{\text{cancels with other side of Schrödinger equation}}+H_1\Psi$$

After cancellations

$$i\hbar\dot{c}_a(t)\psi_a(\mathbf{r})\exp\left(-\tfrac{i}{\hbar}E_a t\right)+i\hbar\dot{c}_b(t)\psi_b(\mathbf{r})\exp\left(-\tfrac{i}{\hbar}E_b t\right)=H_1\Psi \tag{1}$$

Evaluate the inner product of ψ_a and equation (1) to obtain

$$\begin{aligned}&i\hbar\dot{c}_a(t)\exp\left(-\tfrac{i}{\hbar}E_a t\right)\\&\quad=\langle\psi_a|H_1|\Psi\rangle=c_a(t)\langle\psi_a|H_1|\psi_a\rangle\exp\left(-\tfrac{i}{\hbar}E_a t\right)+c_b(t)\langle\psi_a|H_1|\psi_b\rangle\exp\left(-\tfrac{i}{\hbar}E_b t\right)\end{aligned} \tag{2}$$

Evaluate the inner product of ψ_b and equation (1) to obtain

$$\begin{aligned}&i\hbar\dot{c}_b(t)\exp\left(-\tfrac{i}{\hbar}E_b t\right)\\&\quad=\langle\psi_b|H_1|\Psi\rangle=c_a(t)\langle\psi_b|H_1|\psi_a\rangle\exp\left(-\tfrac{i}{\hbar}E_a t\right)+c_b(t)\langle\psi_b|H_1|\psi_b\rangle\exp\left(-\tfrac{i}{\hbar}E_b t\right)\end{aligned} \tag{3}$$

Let it be the case that the following amplitudes vanish.

$$\langle\psi_a|H_1|\psi_a\rangle=0,\quad\langle\psi_b|H_1|\psi_b\rangle=0$$

Then equations (2) and (3) simplify as

$$\begin{aligned} i\hbar\dot{c}_a(t)\exp\left(-\tfrac{i}{\hbar}E_a t\right) &= c_b(t)\langle\psi_a|H_1|\psi_b\rangle\exp\left(-\tfrac{i}{\hbar}E_b t\right) \\ i\hbar\dot{c}_b(t)\exp\left(-\tfrac{i}{\hbar}E_b t\right) &= c_a(t)\langle\psi_b|H_1|\psi_a\rangle\exp\left(-\tfrac{i}{\hbar}E_a t\right) \end{aligned} \tag{4}$$

Let $E_b > E_a$ and let

$$\omega_0 = \frac{E_b - E_a}{\hbar}$$

Rewrite equation (4) as

$$\begin{aligned} \dot{c}_a(t) &= -\frac{i}{\hbar}c_b(t)\langle\psi_a|H_1|\psi_b\rangle\exp(-i\omega_0 t) \\ \dot{c}_b(t) &= -\frac{i}{\hbar}c_a(t)\langle\psi_b|H_1|\psi_a\rangle\exp(i\omega_0 t) \end{aligned}$$

Let the initial conditions be $c_a(0) = 1$ and $c_b(0) = 0$. It was shown in "Perturbation example" that the first-order perturbation solutions are

$$\begin{aligned} &c_a(t) = 1 \\ &c_b(t) = -\frac{i}{\hbar}\int_0^t \langle\psi_b|H_1(\mathbf{r},t')|\psi_a\rangle\exp(i\omega_0 t')\,dt' \end{aligned}$$

The integral is not as bad as it looks because $\psi_a(\mathbf{r})$ and $\psi_b(\mathbf{r})$ are independent of time t. For $H_1(\mathbf{r},t)$ representing an electric plane wave, the integrand reduces to a simple exponential of t' which is easily solved.